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Does God play dice?

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Collaborators

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LI approach - References

Quantum theory as the most robust description of reproducible experiments

Annals of Physics 347 (2014) 45–73

Hans De Raedt[a], Mikhail I. Katsnelson[b], Kristel Michielsen[c,d,*]

Quantum theory as a description of robust experiments: Derivation of the Pauli equation

Annals of Physics 359 (2015) 166–186

Hans De Raedt[a], Mikhail I. Katsnelson[b], Hylke C. Donker[b], Kristel Michielsen[c,d,*]

Logical inference approach to relativistic quantum mechanics: Derivation of the Klein–Gordon equation

Annals of Physics 372 (2016) 74–82

H.C. Donker[a,*], M.I. Katsnelson[a], H. De Raedt[b], K. Michielsen[c]

Logical inference derivation of the quantum theoretical description of Stern–Gerlach and Einstein–Podolsky–Rosen–Bohm experiments

Annals of Physics 396 (2018) 96–118

Hans De Raedt[a], Mikhail I. Katsnelson[b], Kristel Michielsen[c,d,*]

Quantum theory as plausible reasoning applied to data obtained by robust experiments

H. De Raedt[1], M. I. Katsnelson[2] and K. Michielsen[3,4]

PHILOSOPHICAL TRANSACTIONS A

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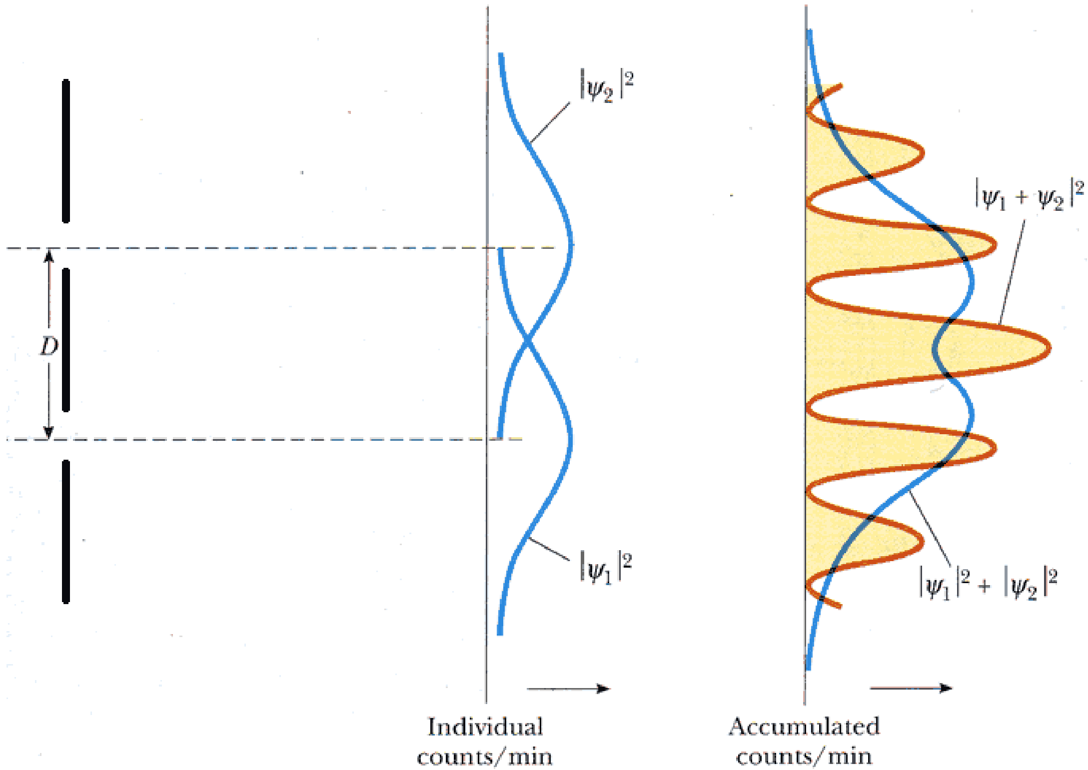
Microworld: waves are corpuscles, corpuscles are waves

Einstein, 1905 – for light (photons)
L. de Broglie, 1924 – electrons and other microparticles

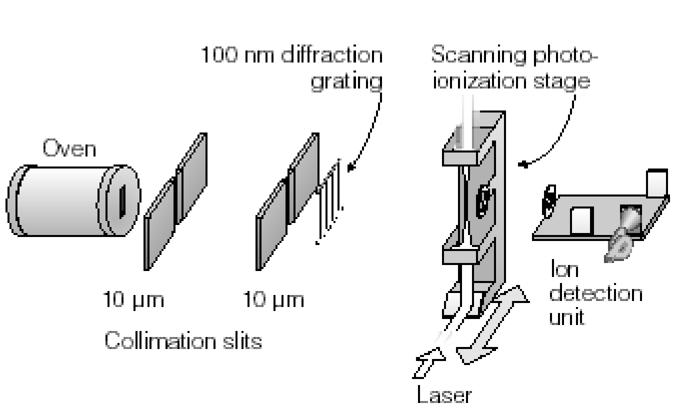


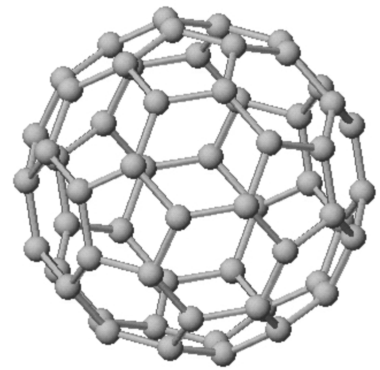
Universal property of matter

Wave–particle duality of C_{60} molecules

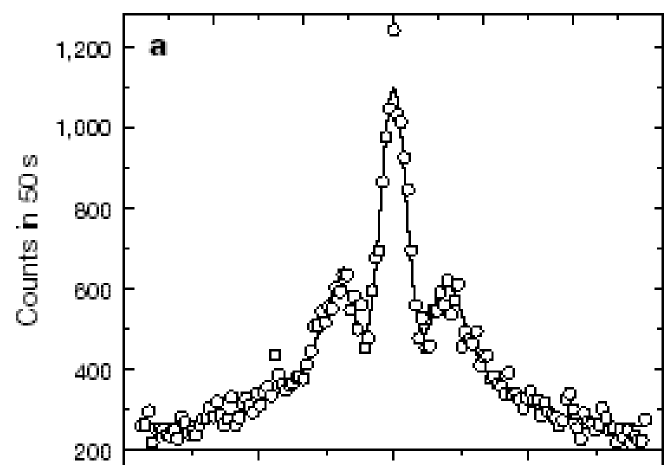
Markus Arndt, Olaf Nairz, Julian Vos-Andreae, Claudia Keller, Gerbrand van der Zouw & Anton Zeilinger

NATURE | VOL 401 | 14 OCTOBER 1999 |



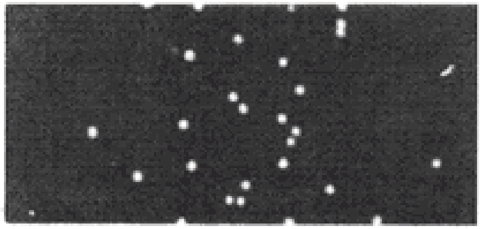


C_{60}

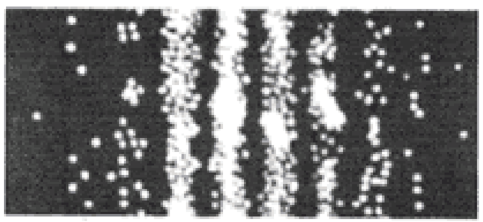


Matter waves for C_{60} molecules

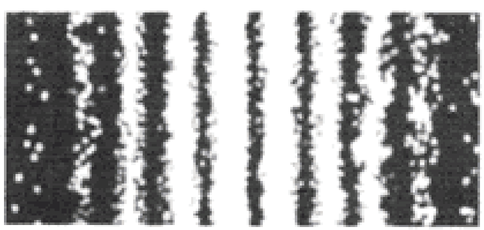
Electrons are particles (you cannot see half of electron) but moves along ***all*** possible directions (interference)



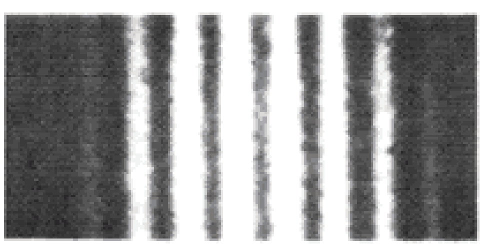
(a) After 28 electrons



(b) After 1000 electrons

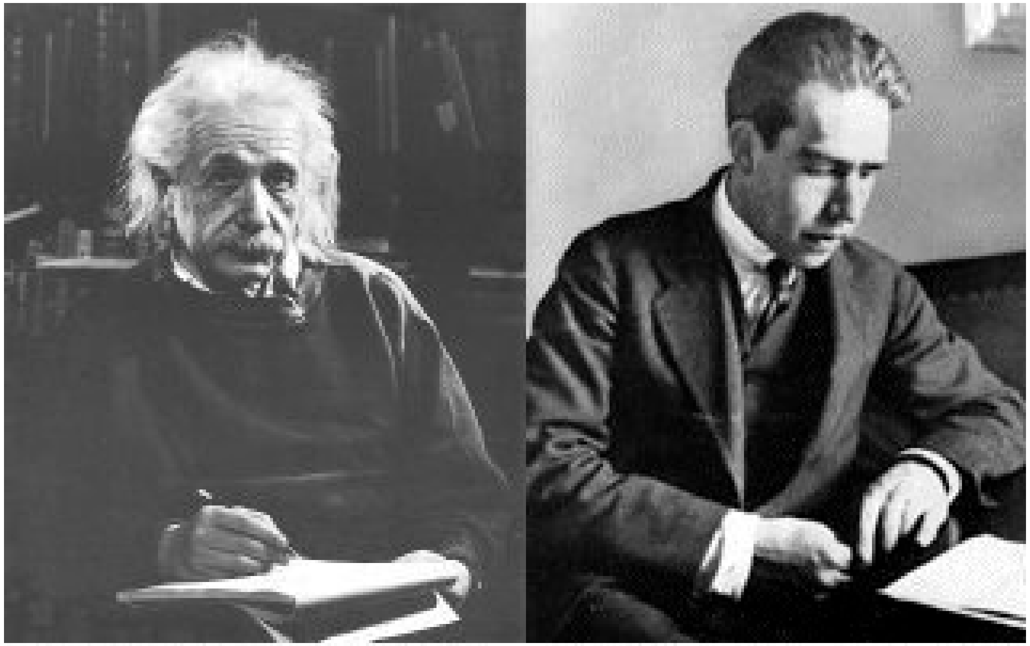


(c) After 10000 electrons



We cannot describe individual events, individual spots seem to be completely random, but ensemble of the spots forms regular interference fridges

Randomness in the foundations of physics?!



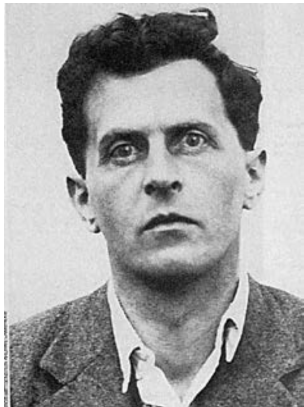
God does not play dice with the universe.
- Albert Einstein

Anyone who is not shocked by Quantum Theory has not understood it. - Niels Bohr

A. Einstein: Quantum mechanics is incomplete; superposition principle does not work in the macroworld

N. Bohr: Classical measurement devices is an important part of quantum reality; we have to describe quantum world in terms of a language created for macroworld

The limits of my language mean the limits of my world
(Ludwig Wittgenstein)

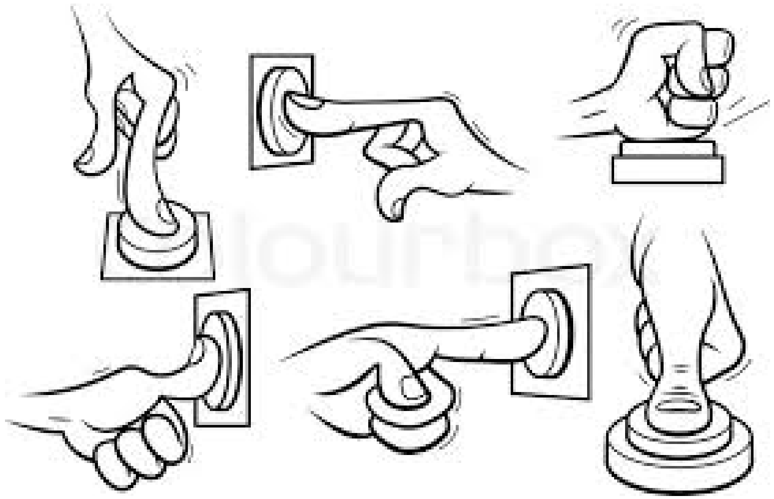


Two ways of thinking

I. Reductionism ("microscopic" approach)

Everything is from water/fire/earth/gauge fields/quantum space-time foam/strings... and the rest is your problem

II. Phenomenology: operating with "black boxes"



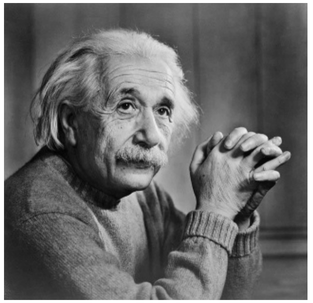
Two ways of thinking II

Knowledge begins, so to speak, in the middle, and leads into the unknown - both when moving upward, and when there is a downward movement. Our goal is to gradually dissipate the darkness in both directions, and the absolute foundation - this huge elephant carrying on his mighty back the tower of truth - it exists only in a fairy tales (Hermann Weyl)

We never know the foundations! How can we have a reliable knowledge without the base?

Is fundamental physics fundamental?

Classical thermodynamics is the only physical theory of universal content which I am convinced will never be overthrown, within the framework of applicability of its basic concepts (A. Einstein)



The laws describing our level of reality are essentially independent on the background laws. I wish our colleagues from *true* theory (strings, quantum gravity, etc....) all kind of success but either they will modify electrodynamics and quantum mechanics at atomic scale (and then they will be wrong) or they will not (and then I do not care). Our way is *down*

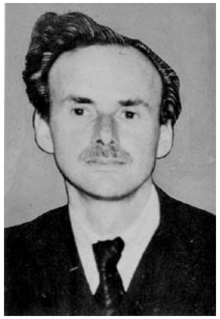
But how can we be sure that we are right?!

Mathematics & Physics

Newton: It is useful to solve (ordinary) differential equations

Maxwell: It is useful to solve *partial* differential equations

Heisenberg, Dirac, von Neumann et al: It is useful to consider state vectors and operators in Hilbert space



But this is much farther from usual human intuition – may be, too far?!

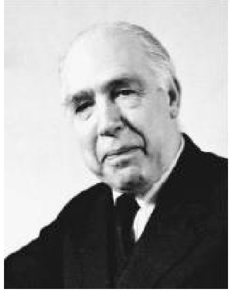
Can we demistify it?!

Unreasonable effectiveness

- Quantum theory describes a vast number of different experiments very well

- WHY ?

- Niels Bohr*:
  It is wrong to think that the task of physics is to find out how nature *is*.
  Physics concerns what we can *say* about nature.



*A. Petersen, "The philosophy of Niels Bohr," Bulletin of the Atomic Scientists 19, 8 – 14 (1963).

Main message of this talk

- Logical inference applied to experiments for which
  1. There is uncertainty about each individual event
  2. The frequencies of observed events are robust with respect to small changes in the conditions

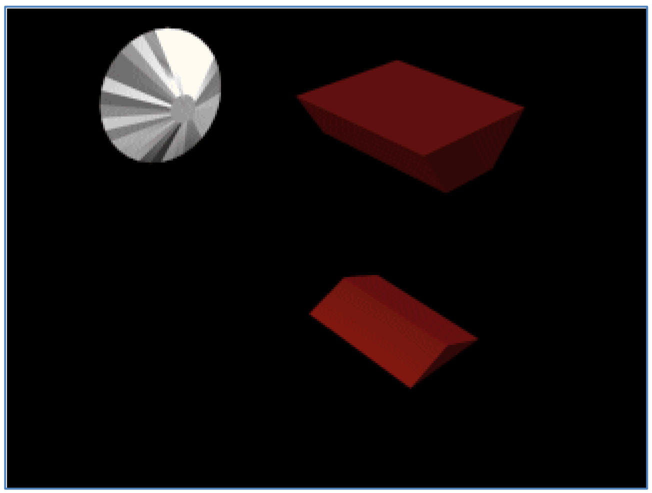
➔ Basic equations of quantum theory

- Not an interpretation of quantum theory

- Derivation based on elementary principles of human reasoning and perception

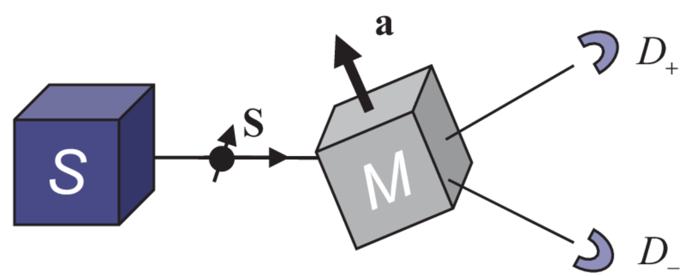
Stern-Gerlach experiment

- Neutral atoms (or neutrons) pass through an inhomogeneous magnetic field



- Inference from the data: directional quantization

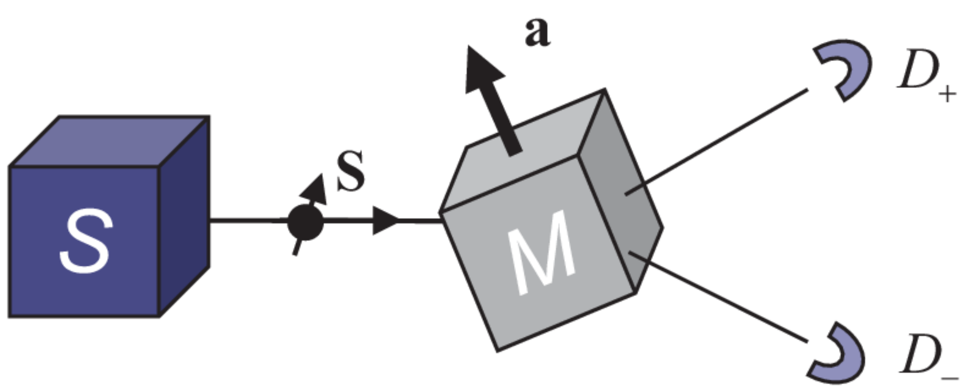
- Idealization



- Source *S* emits particles with magnetic moment
- Magnet *M* sends particle to one of two detectors
- Detectors count every particle

Idealized Stern-Gerlach experiment

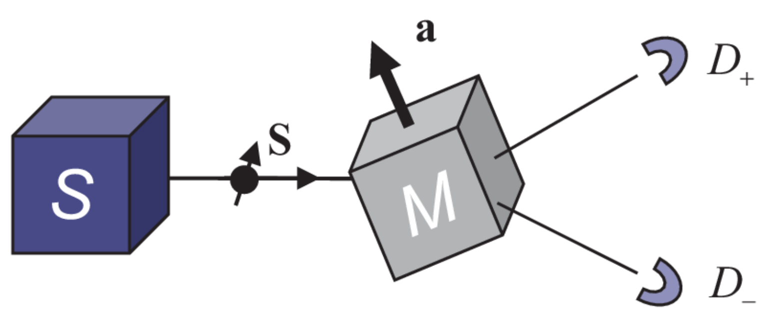
- Event = click of detector D_+ or (exclusive) D_-



- There is uncertainty about each event
  - We do not know how to predict an event with certainty

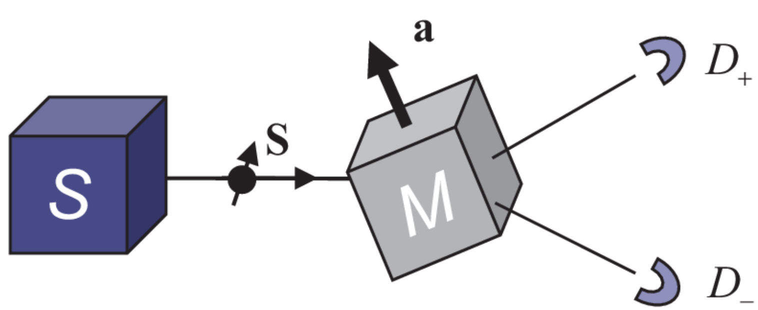
Some reasonable assumptions (1)

- For fixed **a** and fixed source *S*, the frequencies of + and – events are reproducible
- If we rotate the source *S* and the magnet *M* by the same amount, these frequencies do not change



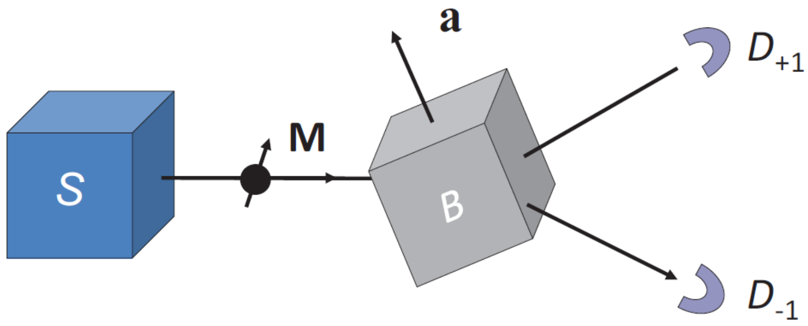
Some reasonable assumptions (2)

- These frequencies are robust with respect to small changes in **a**
- Based on all other events, it is impossible to say with some certainty what the particular event will be (logical independence)



Logical inference

- Shorthand for propositions
  - x=+1 ⇔ D_+ clicks
  - x=-1 ⇔ D_- clicks
  - **M** ⇔the value of **M** is **M**
  - **a** ⇔the value of **a** is **a**
  - **Z** ⇔everything else which is known to be relevant to the experiment but is considered to fixed



- We assign a real number $P(x|\mathbf{M},\mathbf{a},Z)$ between 0 and 1 to express our expectation that detector D_+ or (exclusive) D_- will click and want to derive, not postulate, $P(x|\mathbf{M},\mathbf{a},Z)$ from general principles of rational reasoning
- What are these general principles ?

Plausible, rational reasoning ➔ inductive logic, logical inference

- G. Pólya, R.T. Cox, E.T. Jaynes, …
  - From general considerations about rational reasoning it follows that the plausibility that a proposition *A* (*B*) is true given that proposition *Z* is true may be encoded in real numbers which satisfy

$$\boxed{\begin{array}{l} 0 \le P(A|Z) \le 1 \\ P(A|Z) + P(\bar{A}|Z) = 1 \quad ; \quad \bar{A} = \text{NOT } A \\ P(AB|Z) = P(A|BZ)P(B|Z) \quad ; \quad AB = A \text{ AND } B \end{array}}$$

  - Extension of Boolean logic, applicable to situations in which there is uncertainty about some but not all aspects
    - Kolmogorov's probability theory is an example which complies with the rules of rational reasoning
    - Is quantum theory another example?

Plausible, rational reasoning ➔ logical inference II

- Plausibility
  - Is an intermediate mental construct to carry out inductive logic, rational reasoning, logical inference
  - May express a degree of believe (subjective)
  - May be used to describe phenomena independent of individual subjective judgment

  plausibility ➔ i-prob (inference-probability)

Application to the Stern-Gerlach experiment

We repeat the experiment N times. The number of times that D_+ (D_-) clicks is n_+ (n_-)

i-prob for the individual event is

$$P(x|\mathbf{a}\cdot\mathbf{M},Z) = P(x|\theta,Z) = \frac{1+xE(\theta)}{2} \quad , \quad E(\theta) = E(\mathbf{a}\cdot\mathbf{M},Z) = \sum_{x=\pm1} xP(x|\theta,Z)$$

Dependent on $\cos\theta = \mathbf{a}\cdot\mathbf{M}$ Rotational invariance

Different events are logically independent:

$$P(x_1,\ldots,x_N|\mathbf{a}\cdot\mathbf{M},Z) = \prod_{i=1}^{N} P(x_i|\theta,Z)$$

The i-prob to observe n_+ and n_- events is

$$P(n_{+1},n_{-1}|\theta,N,Z) = N! \prod_{x=\pm1} \frac{P(x|\theta,Z)^{n_x}}{n_x!}$$

How to express robustness?

- Hypothesis H_0: given θ we observe n_+ and n_-
- Hypothesis H_1: given $\theta+\varepsilon$ we observe n_+ and n_-
- The evidence Ev(H_1/H_0) is given by

$$Ev(H_1 \mid H_0) = \ln \frac{P(n_+, n_- \mid \theta+\varepsilon, N, Z)}{P(n_+, n_- \mid \theta, N, Z)} = \sum_{x=\pm 1} n_x \ln \frac{P(x \mid \theta+\varepsilon, Z)}{P(x \mid \theta, Z)} =$$

$$= \sum_{x=\pm 1} n_x \left\{ \varepsilon \frac{P'(x \mid \theta, Z)}{P(x \mid \theta, Z)} - \frac{\varepsilon^2}{2} \left[\frac{P'(x \mid \theta, Z)}{P(x \mid \theta, Z)} \right]^2 + \frac{\varepsilon^2}{2} \frac{P''(x \mid \theta, Z)}{P(x \mid \theta, Z)} \right\} + O(\varepsilon^3)$$

- Frequencies should be robust with respect to small changes in θ ➔ we should minimize, in absolute value, the coefficients of ε, ε^2,...

Remove dependence on ϵ (1)

$$Ev(H_1 | H_0) = \sum_{x=\pm 1} n_x \left\{ \varepsilon \frac{P'(x|\theta,Z)}{P(x|\theta,Z)} - \frac{\varepsilon^2}{2} \left[\frac{P'(x|\theta,Z)}{P(x|\theta,Z)} \right]^2 + \frac{\varepsilon^2}{2} \frac{P''(x|\theta,Z)}{P(x|\theta,Z)} \right\} + O(\varepsilon^3)$$

- Choose $\boxed{P(x|\theta,Z) = \frac{n_x}{N}}$

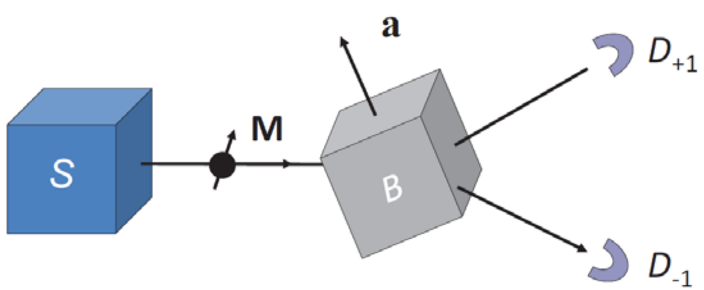
  - Removes the 1st and 3rd term
  - Recover the intuitive procedure of assigning to the i-prob of the individual event, the frequency which maximizes the i-prob to observe the whole data set

Remove dependence on ϵ (2)

$$Ev(H_1 \mid H_0) = \sum_{x=\pm 1} n_x \left\{ \varepsilon \frac{P'(x\mid\theta,Z)}{P(x\mid\theta,Z)} - \frac{\varepsilon^2}{2}\left[\frac{P'(x\mid\theta,Z)}{P(x\mid\theta,Z)}\right]^2 + \frac{\varepsilon^2}{2}\frac{P''(x\mid\theta,Z)}{P(x\mid\theta,Z)} \right\} + O(\varepsilon^3)$$

- Minimizing the 2nd term (Fisher information) for all possible (small) ε and θ

$$\boxed{I_F = \sum_{x=\pm 1} \frac{1}{P(x\mid\theta,Z)}\left(\frac{\partial P(x\mid\theta,Z)}{\partial\theta}\right)^2}$$



$$P(x|\mathbf{a}\cdot\mathbf{M},Z) = P(x|\theta,Z) = \frac{1 \pm x\mathbf{a}\cdot\mathbf{M}}{2}$$

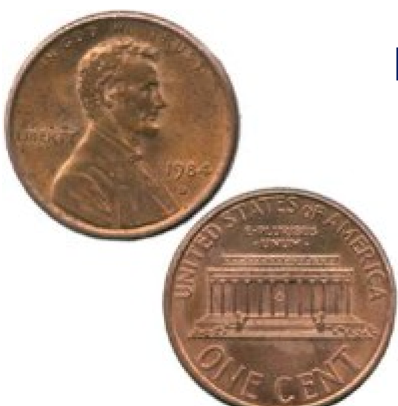
- In agreement with quantum theory of the idealized Stern-Gerlach experiment

Bernoulli trial

Two outcomes (head and tails in coin flypping)

Results are dependent on a single parameter θ which runs a circle (periodicity); what is special in quantum trials?

The results of SG experiment are the most robust, that is, correspond to minimum Fisher information



No assumptions on wave functions, Born rules and other machinery Of quantum physics, just looking for the most robust description of the results of repeating “black box” experiments

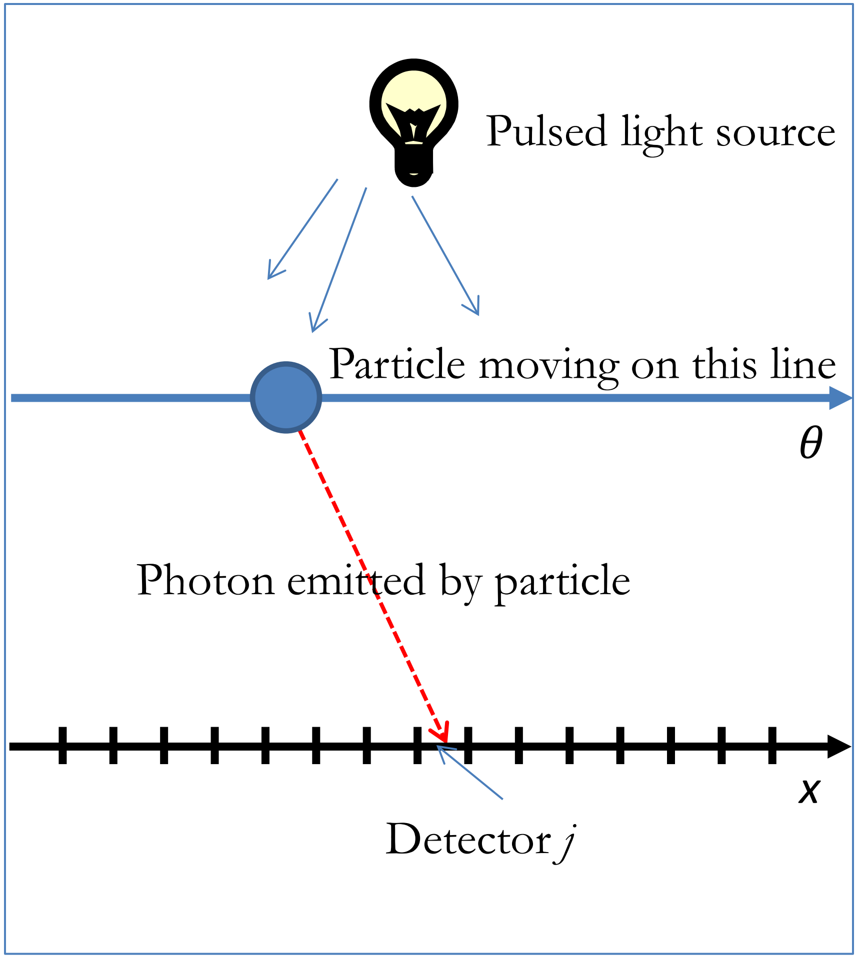
Derivation of basic results of quantum theory by logical inference

- Generic approach
  1. List the features of the experiment that are deemed to be relevant
  2. Introduce the i-prob of individual events
  3. Impose condition of robustness
  4. Minimize functional ➔ equation of quantum theory when applied to experiments in which
     i. There is uncertainty about each event
     ii. The conditions are uncertain
     iii. Frequencies with which events are observed are reproducible and robust against small changes in the conditions

We need to add some “dynamical” information on the system

Logical inference ➔ Schrödinger equation

- Generic procedure:
- Experiment ➔
- The "true" position θ of the particle is uncertain and remains unknown
- i-prob that the particle at unknown position θ activates the detector at position x : $\boxed{P(x\,|\,\theta, Z)}$



Robustness

- Assume that it does not matter if we repeat the experiment somewhere else ➔

$$\boxed{P(x\,|\,\theta,Z)=P(x+\zeta\,|\,\theta+\zeta,Z) \quad ; \quad \zeta \text{ arbitrary}}$$

- Condition for robust frequency distribution ⬄ minimize the functional (Fisher information)

$$\boxed{I_F(\theta)=\int_{-\infty}^{\infty}dx\,\frac{1}{P(x\,|\,\theta,Z)}\left(\frac{\partial P(x\,|\,\theta,Z)}{\partial x}\right)^2}$$

with respect to $P(x|\theta,Z)$

Impose classical mechanics (á la Schrödinger)

- If there is no uncertainty at all ➔classical mechanics ➔ Hamilton-Jacobi equation

$$\boxed{\frac{1}{2m}\left(\frac{\partial S(\theta)}{\partial \theta}\right)^2 + V(\theta) - E = 0} \qquad \text{(X)}$$

- If there is "known" uncertainty

$$\boxed{\int_{-\infty}^{\infty} dx \left[\left(\frac{\partial S(x)}{\partial x}\right)^2 + 2m[V(x) - E]\right] P(x\,|\,\theta, Z) = 0} \qquad \text{(XX)}$$

– Reduces to (X) if $P(x|\theta, Z) \rightarrow \delta(x - \theta)$

Robustness + classical mechanics

- $P(x|\theta, Z)$ can be found by minimizing $I_F(\theta)$ with the constraint that (XX) should hold

➔ We should minimize the functional

$$\boxed{F(\theta) = \int_{-\infty}^{\infty} dx \left\{ \frac{1}{P(x\,|\,\theta, Z)} \left(\frac{\partial P(x\,|\,\theta, Z)}{\partial x} \right)^2 + \lambda \left[\left(\frac{\partial S(x)}{\partial x} \right)^2 + 2m[V(x) - E] \right] P(x\,|\,\theta, Z) \right\}}$$

  - λ = Lagrange multiplier
  - Nonlinear equations for $P(x|\theta, Z)$ and $S(x)$

Robustness + classical mechanics

- Nonlinear equations for $P(x|\theta, Z)$ and $S(x)$ can be turned into linear equations by substituting*

$$\boxed{\psi(x\,|\,\theta,Z)=\sqrt{P(x\,|\,\theta,Z)}e^{iS(x)\sqrt{\lambda}/2}} \quad \Rightarrow$$

$$\boxed{F(\theta)=\int_{-\infty}^{\infty}dx\left\{4\frac{\partial\psi^*(x\,|\,\theta,Z)}{\partial x}\frac{\partial\psi(x\,|\,\theta,Z)}{\partial x}+2m\lambda[V(x)-E]\psi^*(x\,|\,\theta,Z)\psi(x\,|\,\theta,Z)\right\}}$$

- Minimizing with respect to $\psi(x\,|\,\theta,Z)$ yields

$$\boxed{-\frac{\partial^2\psi(x\,|\,\theta,Z)}{\partial x^2}+\frac{m\lambda}{2}\left[V(x)-E\right]\psi(x\,|\,\theta,Z)=0}$$

$\Rightarrow$ Schrödinger equation $\lambda = 4K^{-2} = 4\hbar^{-2}$

*E. Madelung, "Quantentheorie in hydrodynamischer Form," Z. Phys. 40, 322 – 326 (1927)

Time-dependent, multidimensional case

The space is filled by detectors which are fired (or not fired) at some discrete (integer) time $\tau = 1, \ldots, M$

At the very end we have a set of data presented as 0 (no particle in a given box at a given instant or 1

$$\Upsilon = \{\boldsymbol{j}_{n,\tau} | \boldsymbol{j}_{n,\tau} \in [-L^d, L^d];\ n = 1, \ldots, N;\ \tau = 1, \ldots, M\}$$

or, denoting the total counts of voxels $\boldsymbol{j}$ at time τ by $0 \leq k_{\boldsymbol{j},\tau} \leq N$, the experiment produces the data set

$$\mathcal{D} = \left\{ k_{\boldsymbol{j},\tau} \middle| \tau = 1, \ldots, M; N = \sum_{\boldsymbol{j} \in [-L^d, L^d]} k_{\boldsymbol{j},\tau} \right\}. \tag{55}$$

Logical independence of events:

$$P(\mathcal{D}|\theta_1, \ldots, \theta_M, N, Z) = N! \prod_{\tau=1}^{M} \prod_{\boldsymbol{j} \in [-L^d, L^d]} \frac{P(\boldsymbol{j}|\theta_\tau, \tau, Z)^{k_{\boldsymbol{j},\tau}}}{k_{\boldsymbol{j},\tau}!}$$

Time-dependent case II

Homogeneity of the space: $P(\boldsymbol{j}|\boldsymbol{\theta}, Z) = P(\boldsymbol{j}+\boldsymbol{\zeta}|\boldsymbol{\theta}+\boldsymbol{\zeta}, Z)$

Evidence: $$\mathrm{Ev} = \sum_{\boldsymbol{j},\tau} \sum_{i,i'=1}^{d} \frac{\epsilon_{i,\tau}\epsilon_{i',\tau}}{P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)} \frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}{\partial\theta_i} \frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}{\partial\theta_{i'}}$$

$$\mathrm{Ev} = \sum_{\boldsymbol{j},\tau} \left(\sum_{i=1}^{d} \frac{\epsilon_{i},\tau}{\sqrt{P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}} \frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}{\partial\theta_i} \right)^2 \geq 0,$$

and, by using the Cauchy–Schwarz inequality, that

$$\mathrm{Ev} \leq \sum_{\boldsymbol{j},\tau} \left(\sum_{i=1}^{d} \epsilon_{i,\tau}^2 \right) \left(\sum_{i=1}^{d} \frac{1}{P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)} \left(\frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}{\partial\theta_i} \right)^2 \right)$$

$$\widehat{\epsilon}^2 = \max_{i,\tau} \epsilon_{i,\tau}^2$$

$$\leq d\widehat{\epsilon}^2 \sum_{\boldsymbol{j},\tau} \sum_{i=1}^{d} \frac{1}{P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)} \left(\frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau,\tau,Z)}{\partial\theta_i} \right)^2,$$

Time-dependent case III

Minimizing Fisher information: $$I_F = \sum_{\boldsymbol{j},\tau} \sum_{i=1}^{d} \frac{1}{P(\boldsymbol{j}|\boldsymbol{\theta}_\tau, \tau, Z)} \left(\frac{\partial P(\boldsymbol{j}|\boldsymbol{\theta}_\tau, \tau, Z)}{\partial \theta_i} \right)^2$$

Taking into account homogeneity of space; continuum limit:

$$I_F = \int d\boldsymbol{x} \int dt \sum_{i=1}^{d} \frac{1}{P(\boldsymbol{x}|\boldsymbol{\theta}(t), t, Z)} \left(\frac{\partial P(\boldsymbol{x}|\boldsymbol{\theta}(t), t, Z)}{\partial x_i} \right)^2$$

Hamilton – Jacobi equations:

$$\frac{\partial S(\boldsymbol{\theta}, t)}{\partial t} + \frac{1}{2m} \left(\nabla S(\boldsymbol{\theta}, t) - \frac{q}{c} \boldsymbol{A}(\boldsymbol{\theta}, t) \right)^2 + V(\boldsymbol{\theta}, t) = 0$$

Time-dependent case IV

Minimizing functional:

$$F = \int d\mathbf{x} \int dt \sum_{i=1}^{d} \left\{ \frac{1}{P(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)} \left(\frac{\partial P(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)}{\partial x_i} \right)^2 + \lambda \left[\frac{\partial S(\mathbf{x}, t)}{\partial t} + \frac{1}{2m} \left(\frac{\partial S(\mathbf{x}, t)}{\partial x_i} - \frac{q}{c} \mathbf{A}(\mathbf{x}, t) \right)^2 + V(\mathbf{x}, t) \right] P(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \right\}$$

Substitution $\psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) = \sqrt{P(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)} e^{iS(\mathbf{x},t)\sqrt{\lambda}/2}$

Equivalent functional for minimization:

$$\begin{aligned} Q = 2 \int d\mathbf{x} \int dt \, \Bigg\{ & mi\sqrt{\lambda} \Bigg[\psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \frac{\partial \psi^*(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)}{\partial t} \\ & - \psi^*(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \frac{\partial \psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)}{\partial t} \Bigg] \\ & + 2 \sum_{j=1}^{d} \left(\frac{\partial \psi^*(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)}{\partial x_j} + \frac{iq\sqrt{\lambda}}{2c} A_j(\mathbf{x}, t) \psi^*(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \right) \\ & \times \left(\frac{\partial \psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z)}{\partial x_j} - \frac{iq\sqrt{\lambda}}{2c} A_j(\mathbf{x}, t) \psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \right) \\ & + m\lambda V(\mathbf{x}, t) \psi^*(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \psi(\mathbf{x}|\boldsymbol{\theta}(t), t, Z) \Bigg\}, \end{aligned}$$

$$\lambda = 4/\hbar^2$$

Time-dependent case V

Time-dependent Schrödinger equation

$$i\hbar\frac{\partial\psi(\boldsymbol{x}|\boldsymbol{\theta}(t),t,Z)}{\partial t}=\left[-\frac{\hbar^2}{2m}\sum_{j=1}^{d}\left(\frac{\partial}{\partial x_j}-\frac{iq}{\hbar c}\boldsymbol{A}(\boldsymbol{x},t)\right)^2+V(x,t)\right]\psi(\boldsymbol{x}|\boldsymbol{\theta}(t),t,Z)$$

It is linear (superposition principle) which follows from classical Hamiltonian (kinetic energy is $mv^2/2$) and, inportantly, from building one complex function from two real (S and $S+2\pi\hbar$ are equivalent).

A very nontrivial operation dictated just by desire to simplify the problem as much as possible (to pass from nonlinear to linear equation).

Requires further careful thinking!

Separation of conditions principle

H. De Raedt, M. I. Katsnelson, D. Willsch, K. Michielsen, Quantum theory does not need postulates. arXiv:1805.08583

LI allows to derive also Pauli equation, Klein-Gordon equation (Dirac is in progress) but... Superposition principle arises as a trick. Why linear equation? Why wave function? Last not least – what about *open* quantum systtems?

Slightly different view but also based on data analysis

Standard logic: Shrödinger equation ⟶ von Neumann prescription ⟶ description of meaurements. We invert this logic!

Starting point: the way how we deal with the data (reproduced as binary sequences)

Von Neumann theory of measurement (1932)

Density matrix for subsystem A of a total system A + B

$$\rho(\alpha,\alpha') = Tr_\beta \Psi^*(\alpha',\beta)\Psi(\alpha,\beta)$$

$$\rho = \sum_a W_a |a\rangle\langle a|$$

Pure state $\rho = |a\rangle\langle a|$

$\rho^2 = \rho$

Mixed state $Tr\rho^2 < Tr\rho$

Two ways of evolution

1. Unitary evolution

$$i\hbar\frac{\partial\rho}{\partial t} = [H,\rho]$$

$$\rho(t) = \exp(iHt/\hbar)\rho(0)\exp(-iHt/\hbar)$$

Entropy is conserved

$$S = -Tr\rho\ln\rho$$

2. Nonequilibrium evolution by the measurement

$$\rho_{after} = \sum_n P_n \rho_{before} P_n$$

$$P_n = |n\rangle\langle n|$$

$$S_{after} > S_{before}$$

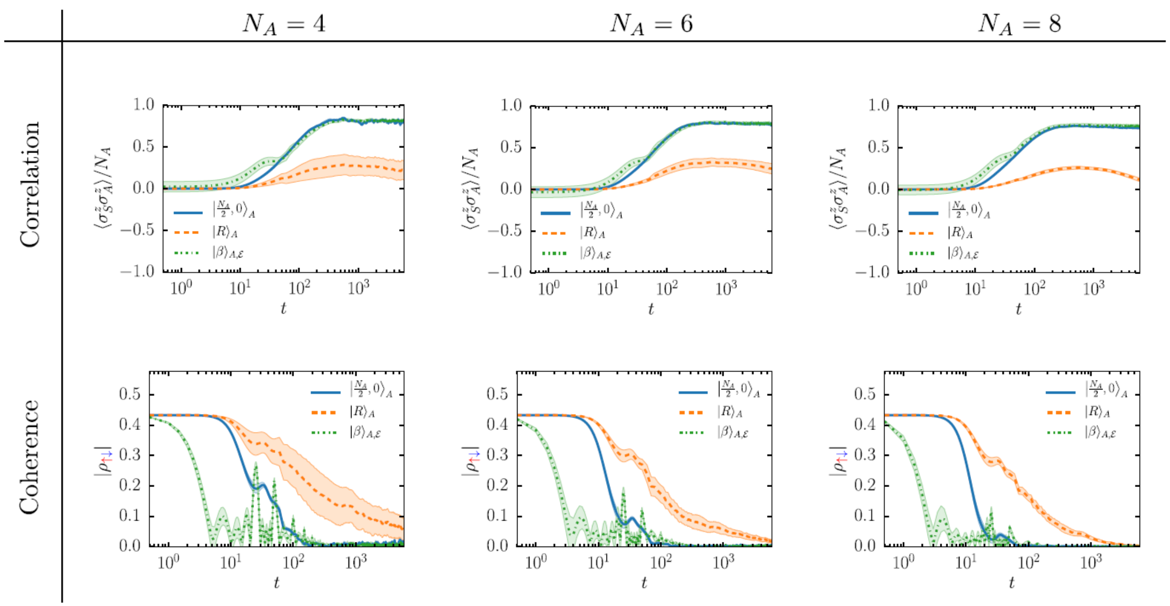
Density matrix after the measurement is diagonal in *n*-representation

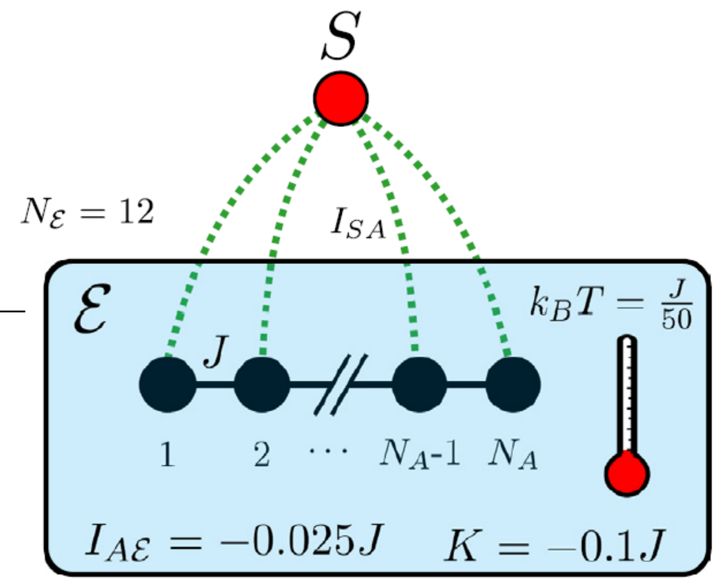
Direct attempts to simulate measurement as interaction with measuring device plus decoherence by environment

Quantum dynamics of a small symmetry breaking measurement device

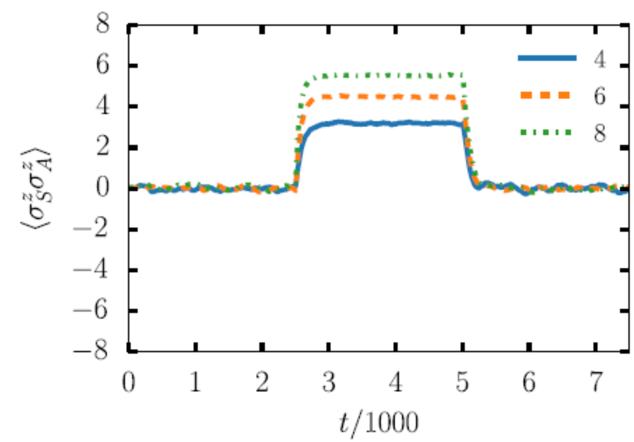
H.C. Donker [a,*], H. De Raedt [b], M.I. Katsnelson [a]

Annals of Physics 396 (2018) 137–146





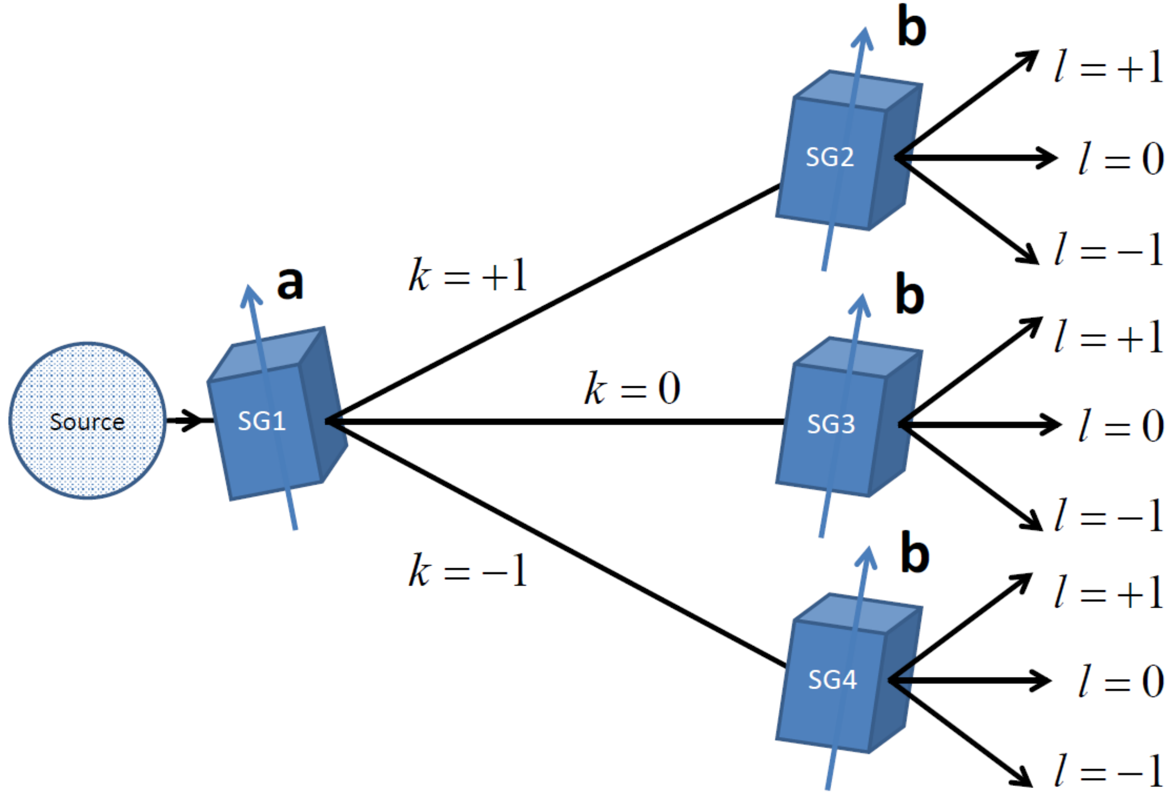
Decay off-diagonal elements of density matrix - yes



Stability test – no stability

Separation procedure

Double SG experiment with three possible outcomes ("spin 1") is generic enough



The first SG device prepares the initial state for the second device

Separation procedure II

The data set for the first device

$$\mathscr{K} = \{k_n \mid k_n \in \{+1, 0, -1\}\ ;\ n = 1, \ldots, N\}$$

$$f(k|\mathbf{a}, P, N) = \frac{1}{N} \sum_{n=1}^{N} \delta_{k,k_n}$$

P properties of the particles emitted by source

Representation in terms of momenta

$$f(k|\mathbf{a}, P, N) = 1 - m_2(\mathbf{a}, P, N) + \frac{m_1(\mathbf{a}, P, N)}{2} k + \frac{3m_2(\mathbf{a}, P, N) - 2}{2} k^2$$

$$m_p(\mathbf{a}, P, N) = \langle k^p \rangle_{\mathbf{a}} = \frac{1}{N} \sum_{n=1}^{N} k_n^p = \sum_{k=+1,0,-1} k^p f(k|\mathbf{a}, P, N) \quad , \quad p = 0, 1, 2$$

Separation procedure III

Let us try to represent the data as strings (sequences)

$$\mathbf{k} = (+1, 0, -1)^T \qquad \mathbf{f} = \left(f(+1|\mathbf{a}, P, N), f(0|\mathbf{a}, P, N), f(-1|\mathbf{a}, P, N)\right)^T$$

$$\langle 1 \rangle_{\mathbf{a}} = (1,1,1) \cdot \mathbf{f} = \mathbf{Tr}\,(1,1,1) \cdot \mathbf{f} = \mathbf{Tr}\,\mathbf{f} \cdot (1,1,1) = \mathbf{Tr} \begin{pmatrix} f(+1|\mathbf{a}, P, N) & 0 & 0 \\ 0 & f(0|\mathbf{a}, P, N) & 0 \\ 0 & 0 & f(-1|\mathbf{a}, P, N) \end{pmatrix}$$

$$\langle k \rangle_{\mathbf{a}} = \mathbf{k}^T \cdot \mathbf{f} = \mathbf{Tr}\,\mathbf{k}^T \cdot \mathbf{f} = \mathbf{Tr}\,\mathbf{f} \cdot \mathbf{k}^T = \mathbf{Tr} \begin{pmatrix} f(+1|\mathbf{a}, P, N) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -f(-1|\mathbf{a}, P, N) \end{pmatrix}$$

$\langle k^2 \rangle_{\mathbf{a}} = \mathbf{Tr}\,\mathbf{f} \cdot (\mathbf{k}^{(2)})^T$ $\qquad$ $\mathbf{k}^{(2)} = (+1, 0, +1)^T$ $\qquad$ is the other vector

Separation procedure IV

But with matrice multiplication rule we need only two matrices

$$\widetilde{\mathbf{K}} = \begin{pmatrix} +1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad \text{and} \quad \widetilde{\mathbf{F}}(\mathbf{a},P,N) = \begin{pmatrix} f(+1|\mathbf{a},P,N) & 0 & 0 \\ 0 & f(0|\mathbf{a},P,N) & 0 \\ 0 & 0 & f(-1|\mathbf{a},P,N) \end{pmatrix}$$

$$\langle k^p \rangle_{\mathbf{a}} = \mathbf{Tr}\,\widetilde{\mathbf{F}}(\mathbf{a},P,N)\widetilde{\mathbf{K}}^p \quad , \quad p = 0,1,2$$

When we rotate the axis of the first SG device and assume rotational invariance (+1 means along the device axis, -1 means opposite, 0 means perpendicular to the axis, for any direction of the axis)

$$\mathbf{K}(\mathbf{a}) = \mathbf{a}\cdot\mathbf{S} \qquad S^x = \frac{1}{\sqrt{2}}\begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \quad , \quad S^y = \frac{1}{\sqrt{2}}\begin{pmatrix} 0 & -i & 0 \\ +i & 0 & -i \\ 0 & +i & 0 \end{pmatrix} \quad , \quad S^z = \begin{pmatrix} +1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

Nothing is quantum yet, except the assumption of *three* outcomes!

Separation procedure V

Introduce projector operator:

$$\mathbf{M}_k(\mathbf{e}_z)\mathbf{M}_l(\mathbf{e}_z) = \delta_{k,l}\mathbf{M}_k(\mathbf{e}_z)$$

$$\mathbf{M}_k(\mathbf{e}_z) = \mathbb{1} - (S^z)^2 + \frac{k}{2}S^z + \frac{k^2}{2}\left[3(S^z)^2 - 2\mathbb{1}\right]$$

$$= \begin{pmatrix} \frac{k^2+k}{2} & 0 & 0 \\ 0 & 1-k^2 & 0 \\ 0 & 0 & \frac{k^2-k}{2} \end{pmatrix} = \begin{cases} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad , \quad k=+1 \\ \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad , \quad k=0 \\ \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad , \quad k=-1 \end{cases}$$

From rotational invariance:

$$\mathbf{M}_k(\mathbf{a}) = 1 - (\mathbf{a}\cdot\mathbf{S})^2 + \frac{k}{2}\mathbf{a}\cdot\mathbf{S} + \frac{k^2}{2}\left[3(\mathbf{a}\cdot\mathbf{S})^2 - 2\mathbb{1}\right]$$

$$f(k|\mathbf{a},P,N) = \mathbf{Tr}\,\mathbf{F}(P,N)\mathbf{M}_k(\mathbf{a}) = \mathbf{Tr}\,\mathbf{M}_k(\mathbf{a})\mathbf{F}(P,N) = \mathbf{Tr}\,\mathbf{M}_k(\mathbf{a})\mathbf{F}(P,N)\mathbf{M}_k(\mathbf{a})$$

Separation procedure VI

As all SG magnets are assumed to be identical, consistency demands that their description should be the same, that is the filtering property of SG2, SG3 and SG4 should be described by $\mathbf{M}_l(\mathbf{b})$.

The first SG device plays the role of the source for the second device etc. – this is the separaction of conditions requirement!

$$\mathscr{D} = \left\{(k_n, l_n) \mid k_n, l_n \in \{+1, 0, -1\}\ ;\ n = 1, \ldots, N\right\}$$

$$f(k|\mathbf{a}, P, N) = \sum_{l=+1,0,-1} f(k, l|\mathbf{a}, \mathbf{b}, P, N)$$

$$f(k, l|\mathbf{a}, \mathbf{b}, P, N) = \mathbf{Tr}\ \mathbf{M}_l(\mathbf{b})\mathbf{M}_k(\mathbf{a})\mathbf{F}(P, N)\mathbf{M}_k(\mathbf{a})\mathbf{M}_l(\mathbf{b})$$

Consequence: $$f(k|\mathbf{a}, P, N) = \sum_{l=+1,0,-1} f(k, l|\mathbf{a}, \mathbf{b}, P, N)$$

Separation procedure VII

Until now *P* (the properties of source) is arbitrary. Illustration:

$$\mathbf{F}(P,N) = \frac{1}{3}\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \qquad f(k|\mathbf{a},P,N) = \mathbf{Tr}\,\mathbf{M}_k(\mathbf{a})\mathbf{F}(P,N)\mathbf{M}_k(\mathbf{a}) = \frac{1}{3}$$

(sourse of unpolarized particles, full isotropy in single SG)

$$\begin{aligned} f(k,l|\mathbf{a},\mathbf{b},P,N) &= \mathbf{Tr}\,\mathbf{M}_l(\mathbf{b})\mathbf{M}_k(\mathbf{a})\mathbf{F}(P,N)\mathbf{M}_k(\mathbf{a})\mathbf{M}_l(\mathbf{b}) \\ &= \begin{cases} \frac{1}{12}(1+\mathbf{a}\cdot\mathbf{b})^2 \quad , \quad k=l=+1,-1 \\ \frac{1}{3}(\mathbf{a}\cdot\mathbf{b})^2 \quad , \quad k=l=0 \\ \frac{1}{12}(1-\mathbf{a}\cdot\mathbf{b})^2 \quad , \quad (k,l)=(+1,-1),(-1,+1) \\ \frac{1}{6}(1-(\mathbf{a}\cdot\mathbf{b})^2) \quad , \quad (k,l)=(+1,0),(-1,0),(0,+1),(0,-1) \end{cases} \end{aligned}$$

This is the result of QM – but strictly speaking not the *derivation*

$$\mathrm{SOC} \models \mathrm{QT}$$

Separation procedure VIII

Dependence on parameters (e.g., time) $\mathscr{D}(\lambda)$ $f(k,l|\mathbf{a},\mathbf{b},P,N,\lambda)$

$$\langle k^p \rangle_\lambda = \mathbf{Tr}\,\mathbf{F}(P,N,\lambda)\mathbf{K}^p(\mathbf{a}) \quad , \quad p = 0,1,2$$

$$\mathbf{Tr}\,\mathbf{F}(P,N,\lambda) = 1 \qquad \mathbf{Tr}\,\frac{\partial^n \mathbf{F}(P,N,\lambda)}{\partial \lambda^n} = 0 \quad , \quad n > 0$$

Traceless matrix is a commutator

K. Shoda, “Einige Sätze über Matrizen,” Jap. J. Math. **13**, 361–365 (1936).
A. A. Albert and B. Muckenhoupt, “On matrices of trace zeros,” Michigan Math. J. , 1–3 (1957).

$$\frac{\partial \mathbf{F}(P,N,\lambda)}{\partial \lambda} = [Y(\lambda), Z(\lambda)]$$

$\mathbf{F}(P,N,\lambda)$ is a Hermitian (non-negative definite) matrix

$\mathbf{F}(P,N,\lambda) = U^\dagger(\lambda)D(\lambda)U(\lambda)$ $D(\lambda)$ is diagonal

Separation procedure IX

$$\frac{\partial \mathbf{F}(P,N,\lambda)}{\partial \lambda} = \left[\mathbf{F}(P,N,\lambda), U^\dagger(\lambda)\frac{\partial U(\lambda)}{\partial \lambda}\right] + U^\dagger(\lambda)\frac{\partial D(\lambda)}{\partial \lambda}U(\lambda)$$

If we assume $\partial D(\lambda)/\partial\lambda = 0$ $\quad \frac{\partial \mathbf{F}(P,N,\lambda)}{\partial \lambda} = i[\mathbf{F}(P,N,\lambda), H(\lambda)]$

$iH(\lambda) = U^\dagger(\lambda)(\partial U(\lambda)/\partial\lambda)$

H is Hermitian and cannot dependent on *F* due to separation requirement

Von Neumann equation:

$$i\hbar\frac{\partial \rho(t)}{\partial t} = [H(t), \rho(t)]$$

If $\rho(t) = |\Psi(t)\rangle\langle\Psi(t)|$ (its eigenvalues are not dependent on time in this case!)

we have Schrödinger equation $i\hbar\frac{\partial}{\partial t}|\Psi(t)\rangle = H(t)|\Psi(t)\rangle$

but to find the “Hamiltonian” one needs other considerations (e.g. like in logical inference part)

To conclude

The way how we deal organize the “data” adds a lot of restrictions on mathematical apparatus which deals with predictions of outcomes of *uncertain* measurements (QT does not predict individual outcomes): (1) Robustness and (2) Separation of conditions

It is not enough to derive QM as a unique theory, some physics should be added but in restricts enormously a class of possible theories

Even if God does not play dice we have to describe the world as if He does

A lot of thing to do but, at least, one can replace (some) (quasi)philosophical declarations by calculations – as we like

Thank you